

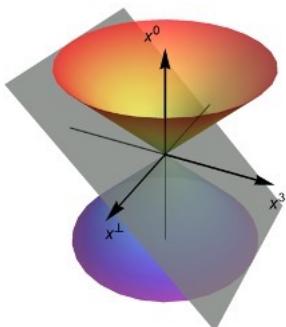
Transverse structure of hadrons beyond leading twist with basis light-front quantization

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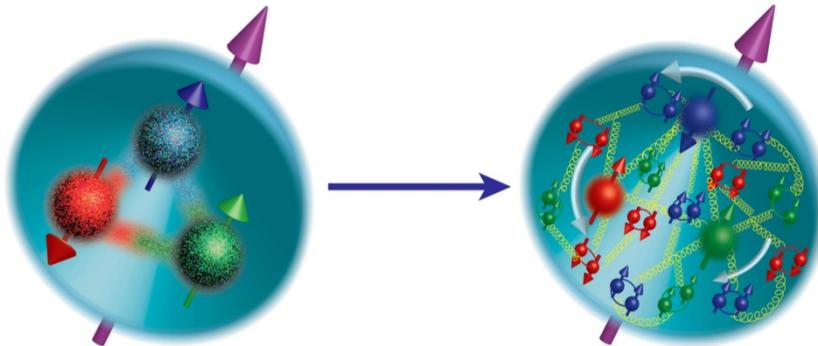
LFQCD Seminars
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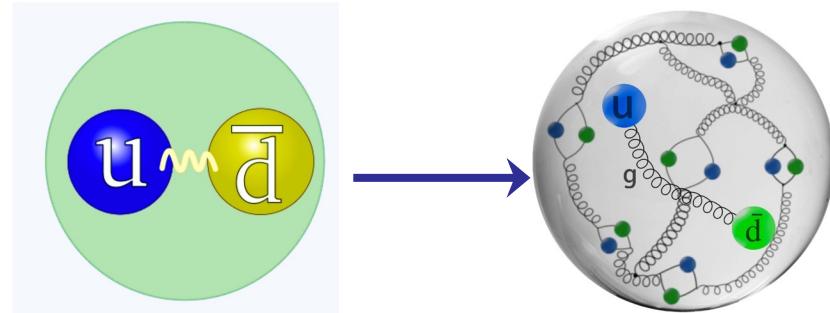
- Hadron structure, PDFs & TMD-PDFs
- Basis light front quantization
- Numerical results

Hadron structures

- Proton is made of
 - 2 up quarks + 1 down quark → valence quarks
 - + any number of quark-antiquark pairs → sea quarks
 - + any number of gluons



$$|\text{proton}\rangle = \psi_{uud} |uud\rangle + \psi_{uudg} |uudg\rangle + \psi_{uudq\bar{q}} |uudq\bar{q}\rangle + \dots$$

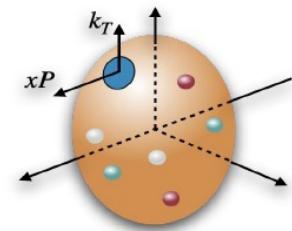


$$|\text{meson}\rangle = \psi_{q\bar{q}} |q\bar{q}\rangle + \psi_{q\bar{q}g} |q\bar{q}g\rangle + \dots$$

Fock sector expansion

- Pion is made of
 - 1 quark + 1 antiquark → valence quarks
 - + any number of quark-antiquark pairs → sea quarks
 - + any number of gluons

✓ Many body system
✓ Rich structure of hadrons



Overview on TMDs (spin 1/2)

Quark correlator

$$\Phi_q^{[\Gamma]} \left(P, S; x = \frac{k^+}{P^+}, \vec{k}_\perp \right) = \frac{1}{2} \int \frac{dz^- dz^\perp}{2(2\pi)^3} e^{ik \cdot z} \langle P, S | \bar{\Psi}_q(0) \Gamma \mathcal{W}(0^\perp, z^\perp) \Psi_q(z) | P, S \rangle \Big|_{z^+ = 0},$$

Parameterization:

8 twist-2 TMDs:

6 T-even terms
2 T-odd terms

$$\Phi^{[\gamma^+]} = f_1 - \frac{\epsilon_\perp^{ij} k_\perp^i S_\perp^j}{M} \textcolor{red}{f_{1T}^\perp},$$

$$\Phi^{[\gamma^+ \gamma^5]} = \Lambda g_{1L} + \frac{k_\perp \cdot \mathbf{S}_\perp}{M} g_{1T},$$

$$\Phi^{[i\sigma^{j+} \gamma^5]} = S_\perp^j h_1 + \Lambda \frac{k_\perp^j}{M} h_{1L}^\perp + S_\perp^i \frac{2k_\perp^i k_\perp^j - (k_\perp)^2 \delta^{ij}}{2M^2} h_{1T}^\perp + \frac{\epsilon_\perp^{ji} k_\perp^i}{M} \textcolor{red}{h_1^\perp},$$

16 twist-3 TMDs:

8 T-even terms
8 T-odd terms

$$\Phi^{[1]} = \frac{M}{P^+} \left[e - \frac{\epsilon_T^{\rho\sigma} k_\perp \rho S_{T\sigma}}{M} \textcolor{red}{e_T^\perp} \right],$$

$$\Phi^{[i\gamma_5]} = \frac{M}{P^+} \left[S_L \textcolor{red}{e_L} - \frac{k_\perp \cdot S_T}{M} \textcolor{red}{e_T} \right],$$

$$\Phi^{[\gamma^\alpha]} = \frac{M}{P^+} \left[-\epsilon_T^{\alpha\rho} S_{T\rho} \textcolor{red}{f_T} - S_L \frac{\epsilon_T^{\alpha\rho} k_\perp \rho}{M} \textcolor{red}{f_L^\perp} - \frac{k_\perp^\alpha k_\perp^\rho - \frac{1}{2} k_\perp^2 g_T^{\alpha\rho}}{M^2} \epsilon_{T\rho\sigma} S_T^\sigma \textcolor{red}{f_T^\perp} + \frac{k_\perp^\alpha}{M} f^\perp \right],$$

$$\Phi^{[\gamma^\alpha \gamma_5]} = \frac{M}{P^+} \left[S_T^\alpha g_T + S_L \frac{k_\perp^\alpha}{M} g_L^\perp - \frac{k_\perp^\alpha k_\perp^\rho - \frac{1}{2} k_\perp^2 g_T^{\alpha\rho}}{M^2} S_{T\rho} g_T^\perp - \frac{\epsilon_T^{\alpha\rho} k_\perp \rho}{M} \textcolor{red}{g_\perp} \right],$$

$$\Phi^{[i\sigma^{\alpha\beta} \gamma_5]} = \frac{M}{P^+} \left[\frac{S_T^\alpha k_\perp^\beta - k_\perp^\alpha S_T^\beta}{M} h_T^\perp - \epsilon_T^{\alpha\beta} \textcolor{red}{h} \right],$$

$$\Phi^{[i\sigma^{+-} \gamma_5]} = \frac{M}{P^+} \left[S_L h_L - \frac{k_\perp \cdot S_T}{M} h_T \right].$$

Jaffe-Ji notation:

f, e → unpolarized quarks

g → longitudinally polarized quarks

h → transversely polarized quarks

1 → the leading twist

L → longitudinally polarized hadron

T → transversely polarized hadron

\perp → existing k_\perp with a non-contracted index

Twist-2 TMDs

$f \rightarrow$ unpolarized quarks

$g \rightarrow$ longitudinally polarized quarks

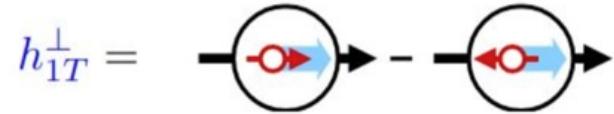
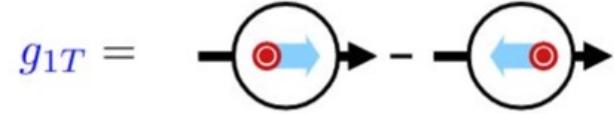
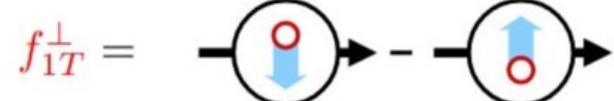
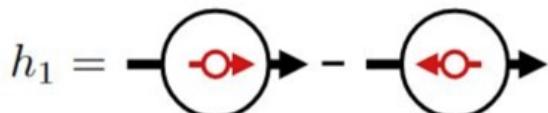
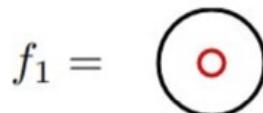
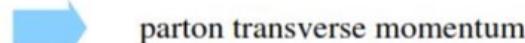
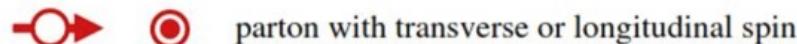
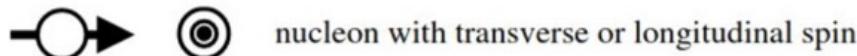
$h \rightarrow$ transversely polarized quarks

$1 \rightarrow$ the leading twist

$L \rightarrow$ longitudinally polarized hadron

$T \rightarrow$ transversely polarized hadron

$\perp \rightarrow$ existing k_\perp with a non-contracted index



Hadron structure: encoded in TMDs

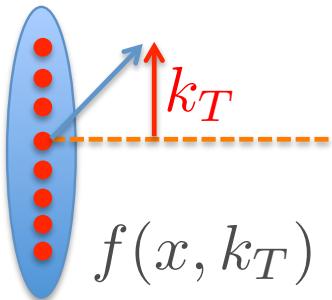
TMD-PDFs

Parton Distribution Functions (PDFs):

Transverse-momentum-dependent PDFs:

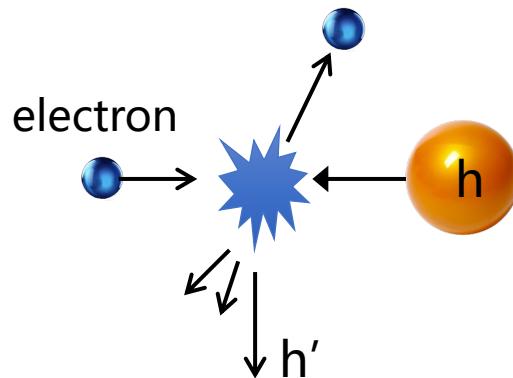
1D longitudinal

1D longitudinal + 2D transverse

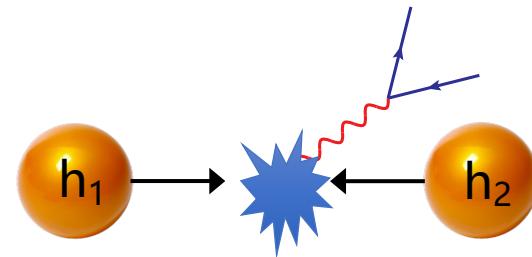


Longitudinal + transverse motion

Probability for finding a parton in a hadron with momentum fraction x and transverse momentum k_T



Semi-inclusive Deep Inelastic Scattering (SIDIS)



Drell-Yan process

SIDIS and Drell-Yan processes: extracting TMDs

Semi-inclusive DIS

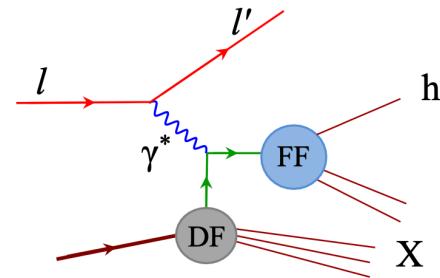
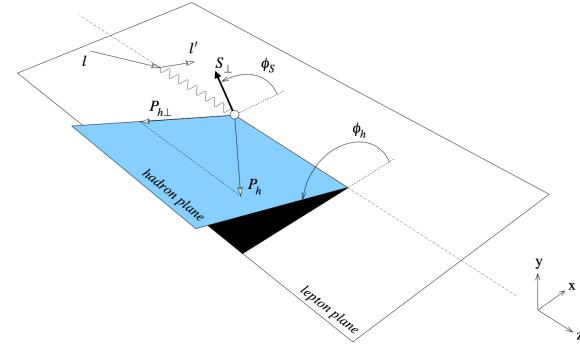
$$\frac{d\sigma}{dxdydzdP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$\left. \begin{aligned} & 1 + \cos \phi_h \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \phi_h} \right) + \cos 2\phi_h \left(\varepsilon A_{UU}^{\cos 2\phi_h} \right) \\ & + \lambda \sin \phi_h \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LU}^{\sin \phi_h} \right) \\ & + S_L \left[\sin \phi_h \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UL}^{\sin \phi_h} \right) + \sin 2\phi_h \left(\varepsilon A_{UL}^{\sin 2\phi_h} \right) \right] \\ & + S_L \lambda \left[\sqrt{1-\varepsilon^2} A_{LL} + \cos \phi_h \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LL}^{\cos \phi_h} \right) \right] \end{aligned} \right]$$

$$\left. \begin{aligned} & \sin(\phi_h - \phi_S) \left(A_{UT}^{\sin(\phi_h - \phi_S)} \right) \\ & + \sin(\phi_h + \phi_S) \left(\varepsilon A_{UT}^{\sin(\phi_h + \phi_S)} \right) \\ & + \sin(3\phi_h - \phi_S) \left(\varepsilon A_{UT}^{\sin(3\phi_h - \phi_S)} \right) \\ & + \sin \phi_S \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin \phi_S} \right) \\ & + \sin(2\phi_h - \phi_S) \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin(2\phi_h - \phi_S)} \right) \end{aligned} \right] \quad \text{SSA} \uparrow$$

$$\left. \begin{aligned} & \cos(\phi_h - \phi_S) \left(\sqrt{(1-\varepsilon^2)} A_{LT}^{\cos(\phi_h - \phi_S)} \right) \\ & + \cos \phi_S \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos \phi_S} \right) \\ & + \cos(2\phi_h - \phi_S) \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos(2\phi_h - \phi_S)} \right) \end{aligned} \right] \quad \text{DSA} \downarrow$$

$$+ S_T \quad \quad \quad + S_T \lambda$$



$$A_{UT}^{\sin(\phi_h - \phi_S)} \propto f_{1T}^\perp \otimes D_1$$

$$A_{UT}^{\sin(\phi_h + \phi_S)} \propto h_1 \otimes H_1^\perp$$

$$A_{UT}^{\sin(3\phi_h - \phi_S)} \propto h_{1T}^\perp \otimes H_1^\perp$$

$$A_{UT}^{\sin(\phi_S)} \propto \frac{M}{Q} (f_T \otimes D_1 + h_1 \otimes H_1^\perp + \dots)$$

$$A_{UT}^{\sin(2\phi_h - \phi_S)} \propto \frac{M}{Q} (h_T \otimes H_1^\perp + h_T^\perp H_1^\perp + \dots)$$

$$A_{LT}^{\cos(\phi_h - \phi_S)} \propto g_{1T} \otimes D_1$$

$$A_{LT}^{\cos(\phi_S)} \propto \frac{M}{Q} (g_T \otimes D_1 + e_T \otimes H_1^\perp + \dots)$$

$$A_{LT}^{\cos(2\phi_h - \phi_S)} \propto \frac{M}{Q} (e_T \otimes H_1^\perp + e_T^\perp \otimes H_1^\perp + \dots)$$

Twist-2

Twist-3

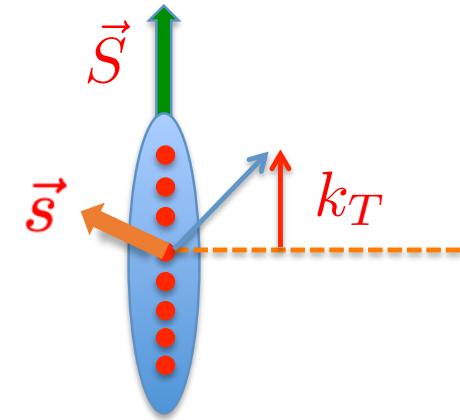
[Bacchetta, et al, JHEP 02 (2007) 093⁷

Quark TMD-PDF for hadrons

Quark correlator

$$\Phi_q^{[\Gamma]} \left(P, S; x = \frac{k^+}{P^+}, k_\perp \right) = \frac{1}{2} \int \frac{dz^- dz^\perp}{2(2\pi)^3} e^{ik \cdot z} \langle P, S | \bar{\Psi}_q(0) \Gamma \boxed{\mathcal{W}(0^\perp, z^\perp)} \Psi_q(z) | P, S \rangle \Big|_{z^+ = 0},$$

$\mathcal{W} \rightarrow 1$



Twist-2: $\Phi^{[\gamma^+]}(x, k_\perp) = f_1(x, k_\perp),$

.....

(averaging hadron spin S)

Twist-3: $\Phi^{[1]}(x, k_\perp) = \frac{M}{P^+} e(x, k_\perp),$

.....

Hadron states:

$$|\text{proton}\rangle = \underbrace{\psi_{uud}}_{|uud\rangle} |uud\rangle + \underbrace{\psi_{uudg}}_{|uudg\rangle} |uudg\rangle + \dots$$

$$|\text{meson}\rangle = \underbrace{\psi_{q\bar{q}}}_{|q\bar{q}\rangle} |q\bar{q}\rangle + \underbrace{\psi_{q\bar{q}g}}_{|q\bar{q}g\rangle} |q\bar{q}g\rangle + \dots$$

Light-front wave functions

Overlap form of TMDs:

Proton: $f_1 \sim \langle P, S | \bar{\Psi} \gamma^+ \Psi | P, S \rangle \sim \int [D] \psi_{uud}^* \psi_{uud} + \int [D] \psi_{uudg}^* \psi_{uudg}$

Meson: $f_1 \sim \langle P | \bar{\Psi} \gamma^+ \Psi | P \rangle \sim \int [D] \psi_{q\bar{q}}^* \psi_{q\bar{q}} + \int [D] \psi_{q\bar{q}g}^* \psi_{q\bar{q}g}$

Equation of motion relation

Using the unintegrated quark correlator

$$\Phi_{\alpha\beta}(x, k_\perp) = \int \frac{d^4 z}{(2\pi)^4} e^{ik \cdot z} \langle P | \bar{\psi}_\beta(0) \psi_\alpha(z) | P \rangle ,$$

Inserting the equation of motion

$$0 = \int \frac{d^4 z}{(2\pi)^4} e^{ik \cdot z} \langle P | \bar{\psi}(0) \Gamma(iD(z) - m) \psi(z) | P \rangle ,$$

Through some algebra,

$$0 = \int \frac{dz^- d^2 z_\perp}{2(2\pi)^3} e^{ik \cdot z} \langle P | \bar{\psi}(0) [k^+ \mathbb{1} - m \gamma^+ + g A^j(0) i \sigma^{j+}] \psi(z) | P \rangle|_{z^+=0}.$$

$$(2\pi)^3 \cdot \langle \bar{q} q \rangle [x^2 - \dots] + g_{11} \langle \bar{q} q \rangle \langle \bar{q} q \rangle - \gamma_1$$


 $e(x, k_\perp) = \frac{m}{M} \frac{f_1(x, k_\perp)}{x} + \tilde{e}(x, k_\perp),$
involve qqq interactions

twist-3	twist-2	genuine twist-3
---------	---------	-----------------

$$|\text{proton}\rangle = \psi_{uud}|uud\rangle + \psi_{uudg}|uudg\rangle + \dots$$

$$\text{twist-2 TMD: } f_1(x, k_\perp) \sim \langle P | \bar{\psi} \gamma^+ \psi | P \rangle \sim \int [D] \psi_{uud}^* \psi_{uud} + \int [D] \psi_{uudg}^* \psi_{uudg}$$

genuine twist-3 TMD: $\tilde{e}(x, k_\perp) \sim \langle P | \bar{\psi} \sigma^{j+} A_j \psi | P \rangle \sim \int [D] \left[\psi_{uud}^* \psi_{uudg} + \text{h.c.} \right]$

Equation of motion relation [Bacchetta, JHEP 02 (2007) 093]

Quark-quark-gluon correlator

$$\Phi_A^{\alpha[\Gamma]}(x, k_\perp) = g \int \frac{d\xi^- d^2\xi_\perp}{2(2\pi)^3} e^{ik \cdot \xi} \langle P, S | \bar{\psi}(0) \Gamma A^\alpha(\xi) \psi(\xi) | P, S \rangle|_{\xi^+=0}$$

16 terms Genuine twist-3 TMDs:

8 T-even terms

8 T-odd terms

$$\begin{aligned} \frac{1}{2Mx} \text{Tr} [\tilde{\Phi}_{A\alpha} \sigma^{\alpha+}] &= \tilde{h} + i\tilde{e} + \frac{\epsilon_T^{\rho\sigma} k_{\perp\rho} S_{T\sigma}}{M} (\tilde{h}_T^\perp - i\tilde{e}_T^\perp), \\ \frac{1}{2Mx} \text{Tr} [\tilde{\Phi}_{A\alpha} i\sigma^{\alpha+} \gamma_5] &= S_L (\tilde{h}_L + i\tilde{e}_L) - \frac{k_\perp \cdot S_T}{M} (\tilde{h}_T + i\tilde{e}_T), \\ \frac{1}{2Mx} \text{Tr} [\tilde{\Phi}_{A\rho} (g_T^{\alpha\rho} + i\epsilon_T^{\alpha\rho} \gamma_5) \gamma^+] &= \frac{k_\perp^\alpha}{M} (\tilde{f}^\perp - i\tilde{g}^\perp) - \epsilon_T^{\alpha\rho} S_{T\rho} (\tilde{f}_T + i\tilde{g}_T) \\ &- S_L \frac{\epsilon_T^{\alpha\rho} k_{\perp\rho}}{M} (\tilde{f}_L^\perp + i\tilde{g}_L^\perp) - \frac{k_\perp^\alpha k_\perp^\rho - \frac{1}{2} k_\perp^2 g_T^{\alpha\rho}}{M^2} \epsilon_{T\rho\sigma} S_T^\sigma (\tilde{f}_T^\perp + i\tilde{g}_T^\perp), \end{aligned}$$

8 EOM relations:

$$\begin{aligned} e(x, k_\perp^2) &= \frac{m}{M} \frac{f_1(x, k_\perp^2)}{x} + \tilde{e}(x, k_\perp^2), \\ f^\perp(x, k_\perp^2) &= \frac{f_1(x, k_\perp^2)}{x} + \tilde{f}^\perp(x, k_\perp^2), \\ g'_T(x, k_\perp^2) &= \frac{m}{M} \frac{h_{1T}(x, k_\perp^2)}{x} + \tilde{g}'_T(x, k_\perp^2), \\ g_L^\perp(x, k_\perp^2) &= \frac{g_{1L}(x, k_\perp^2)}{x} + \frac{m}{M} \frac{h_{1L}^\perp(x, k_\perp^2)}{x} + \tilde{g}_L^\perp(x, k_\perp^2), \\ g_T^\perp(x, k_\perp^2) &= \frac{g_{1T}(x, k_\perp^2)}{x} + \frac{m}{M} \frac{h_{1T}^\perp(x, k_\perp^2)}{x} + \tilde{g}_T^\perp(x, k_\perp^2), \\ h_T^\perp(x, k_\perp^2) &= \frac{h_{1T}(x, k_\perp^2)}{x} + \tilde{h}_T^\perp(x, k_\perp^2), \\ h_L(x, k_\perp^2) &= \frac{m}{M} \frac{g_{1L}(x, k_\perp^2)}{x} - \frac{k_\perp^2}{M^2} \frac{h_{1L}^\perp(x, k_\perp^2)}{x} + \tilde{h}_L(x, k_\perp^2), \\ h_T(x, k_\perp^2) &= \frac{m}{M} \frac{g_{1T}(x, k_\perp^2)}{x} - \frac{h_{1T}(x, k_\perp^2)}{x} - \frac{k_\perp^2}{M^2} \frac{h_{1T}^\perp(x, k_\perp^2)}{x} + \tilde{h}_T(x, k_\perp^2), \end{aligned}$$

Basis light front quantization (BLFQ)

Light-front Schrödinger equation

[Vary et al, 2008]

$$P^+ P^- |\Psi\rangle = M^2 |\Psi\rangle$$

Convert to

Eigenvalue problem

$$P^- |\beta\rangle = P_\beta^- |\beta\rangle$$

- P^- : light-front Hamiltonian
- $|\beta\rangle$: mass eigenstate
- P_β^- : eigenvalue for $|\beta\rangle$

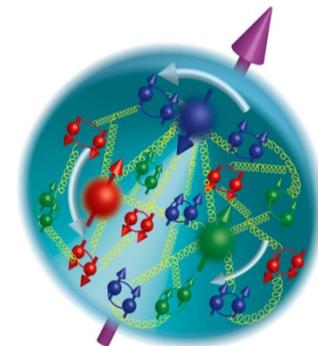
- Fock sector expansion

$$|\text{proton}\rangle = \psi_{uud} |uud\rangle + \psi_{uudg} |uudg\rangle + \dots$$

- Transverse: 2D harmonic oscillator basis: $\Phi_{n,m}^b(\vec{p}_\perp)$.
- Longitudinal: plane-wave basis, labeled by k .
- Basis truncation:

$$\begin{aligned} \sum_i (2n_i + |m_i| + 1) &\leq N_{max}, \\ \sum_i k_i &= K_{max}. \end{aligned}$$

N_{max}, K_{max} are basis truncation parameters.



Proton with one dynamic gluon

$$P^+ P^- |\Psi\rangle = M^2 |\Psi\rangle$$

$$|\text{proton}\rangle = \psi_{uud} |uud\rangle + \psi_{uudg} |uudg\rangle$$

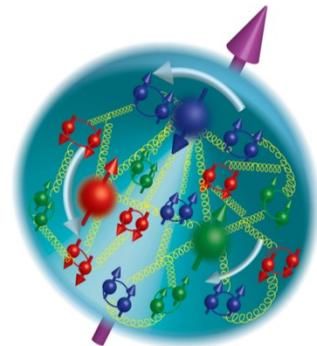
QCD Interaction:

$$\begin{aligned} P_{\text{QCD}}^- = & \int dx^- d^2x^\perp \left\{ \frac{1}{2} \bar{\psi} \gamma^+ \frac{m_0^2 + (i\partial^\perp)^2}{i\partial^+} \psi \right. \\ & - \frac{1}{2} A_a^i [m_g^2 + (i\partial^\perp)^2] A_a^i + g_s \bar{\psi} \gamma_\mu T^a A_a^\mu \psi \\ & \left. + \frac{1}{2} g_s^2 \bar{\psi} \gamma^+ T^a \psi \frac{1}{(i\partial^+)^2} \bar{\psi} \gamma^+ T^a \psi \right\}, \end{aligned}$$

[Siqi XU et al, arXiv: 2209.0858]

Confinement:

$$P_C^- P^+ = \frac{\kappa^4}{2} \sum_{i \neq j} \left\{ \{ \vec{r}_{ij}^2 - \frac{\partial_{x_i} (x_i x_j \partial_{x_j})}{(m_i + m_j)^2} \} \right.$$



Parameters:

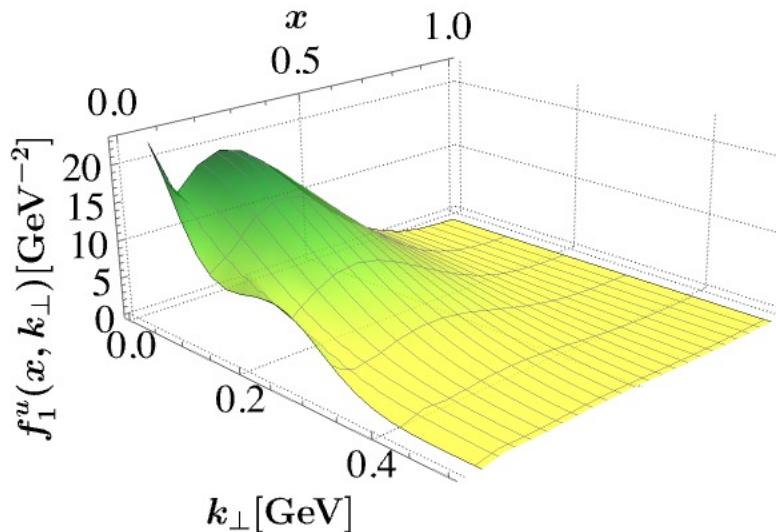
Truncation: Nmax=9, K=16.5

HO parameters: b=0.7GeV, b_{inst}=3GeV

m _u	m _d	m _g	κ	m _f	g
0.31GeV	0.25GeV	0.50GeV	0.54GeV	1.80GeV	2.40

Twist-2 TMDs of proton

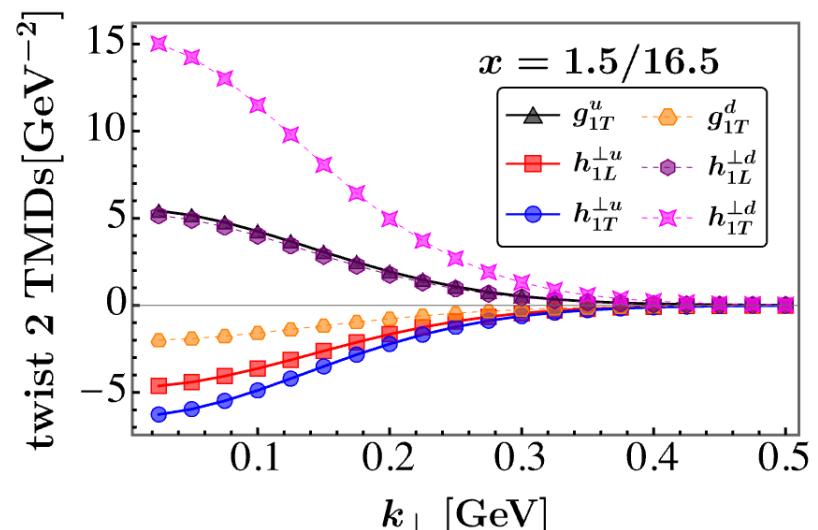
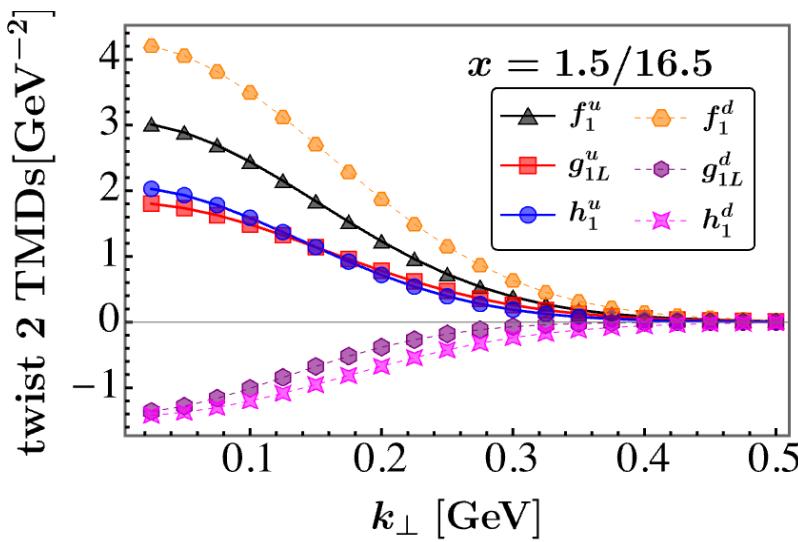
[Hongyao YU et al, in preparation]



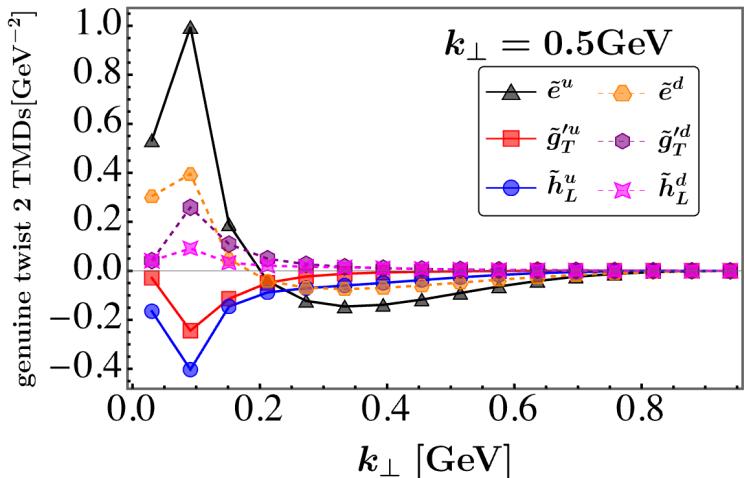
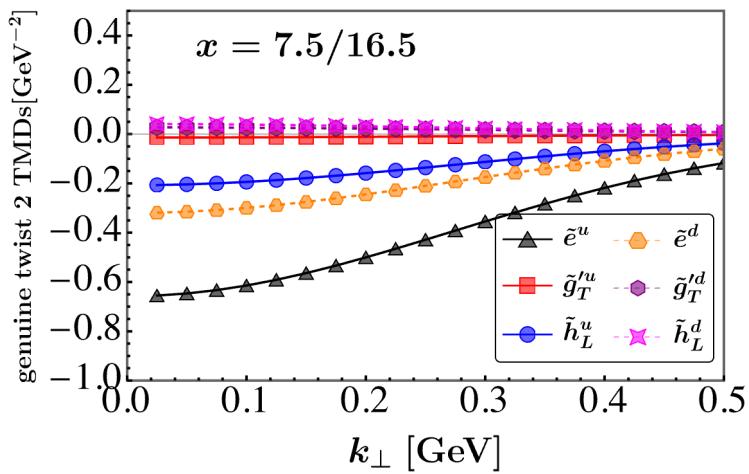
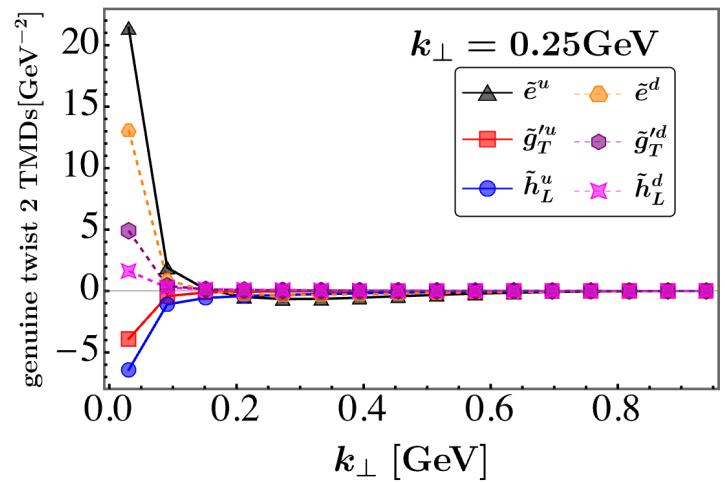
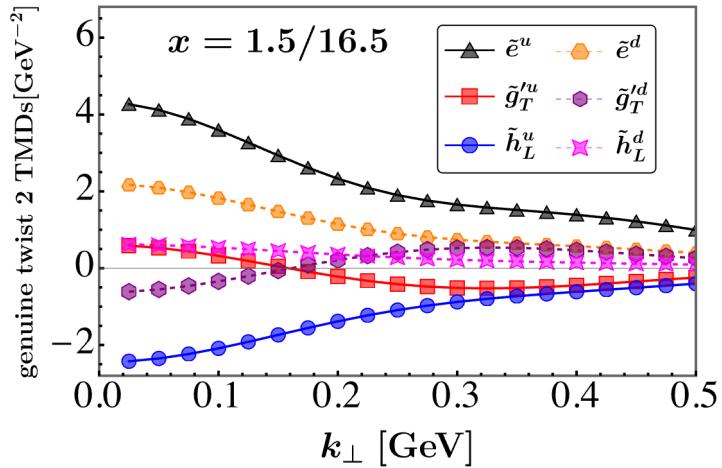
$$\Phi[\gamma^+] = f_1 - \frac{\epsilon_\perp^{ij} k_\perp^i S_\perp^j}{M} \mathbf{f}_{1T}^\perp,$$

$$\Phi[\gamma^+ \gamma^5] = \Lambda g_{1L} + \frac{k_\perp \cdot \mathbf{S}_\perp}{M} g_{1T},$$

$$\Phi[i\sigma^{j+}\gamma^5] = S_\perp^j h_1 + \Lambda \frac{k_\perp^j}{M} h_{1L}^\perp + S_\perp^i \frac{2k_\perp^i k_T^j - (k_\perp)^2 \delta^{ij}}{2M^2} h_{1T}^\perp + \frac{\epsilon_\perp^{ji} k_\perp^i}{M} \mathbf{h}_1^\perp,$$



Genuine twist-3 TMDs of proton

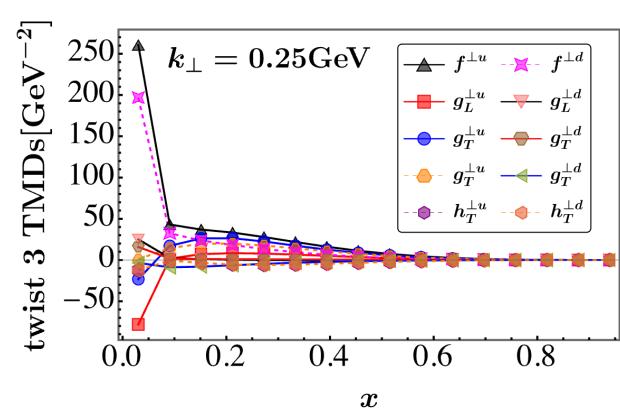
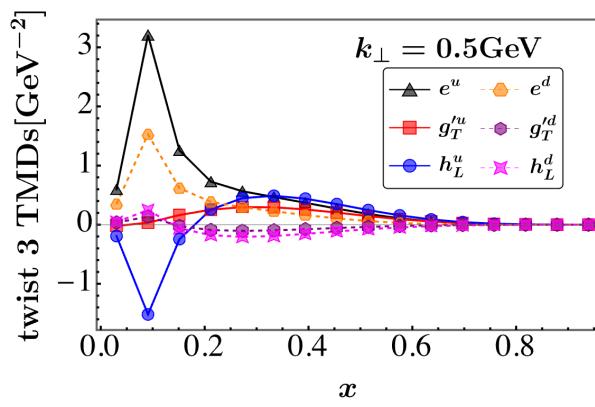
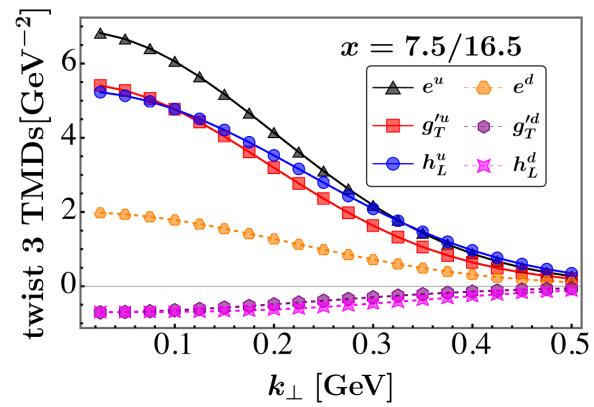
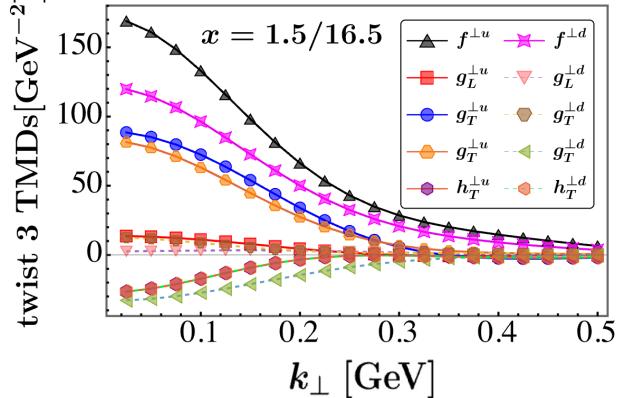
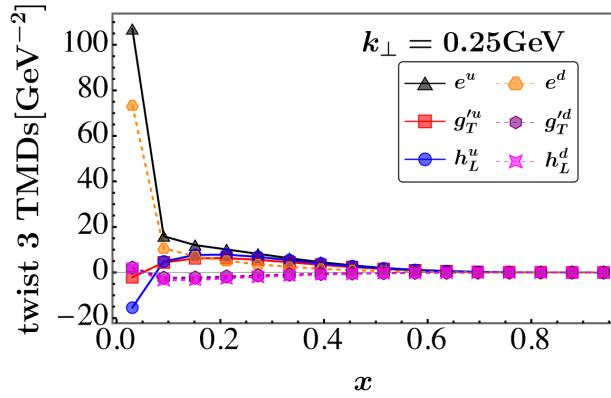
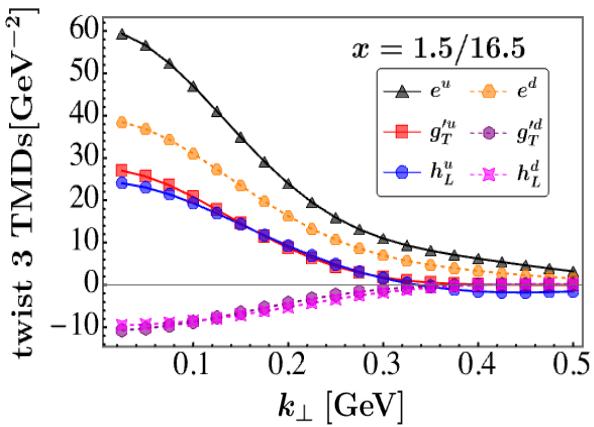


- more concentrating in small x and k_T
- same magnitude as twist-2 parts

$$e(x, k_{\perp}^2) = \frac{m}{M} \frac{f_1(x, k_{\perp}^2)}{x} + \tilde{e}(x, k_{\perp}^2)$$

Twist-3 TMDs of proton

[Zhimin ZHU et al, in preparation]



- more concentrating in small x and k_T
- large than twist-2 TMDs

$$e(x, k_\perp^2) = \frac{m}{M} \frac{f_1(x, k_\perp^2)}{x} + \tilde{e}(x, k_\perp^2)$$

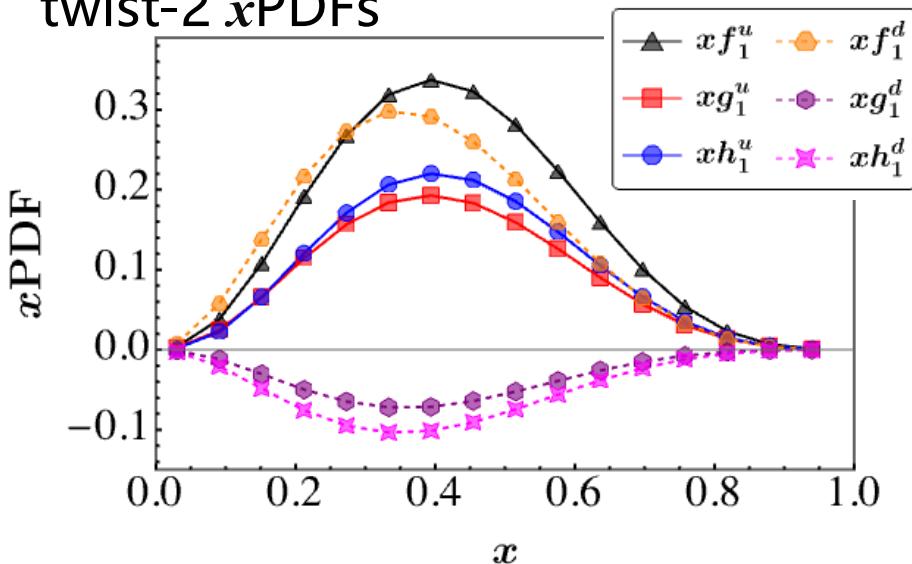
$$g'_T(x, k_\perp^2) = \frac{m}{M} \frac{h_{1T}(x, k_\perp^2)}{x} + \tilde{g}'_T(x, k_\perp^2)$$

$$h_L(x, k_\perp^2) = \frac{m}{M} \frac{g_{1L}(x, k_\perp^2)}{x} - \frac{k_\perp^2}{M^2} \frac{h_{1L}^\perp(x, k_\perp^2)}{x} + \tilde{h}_L(x, k_\perp^2)$$

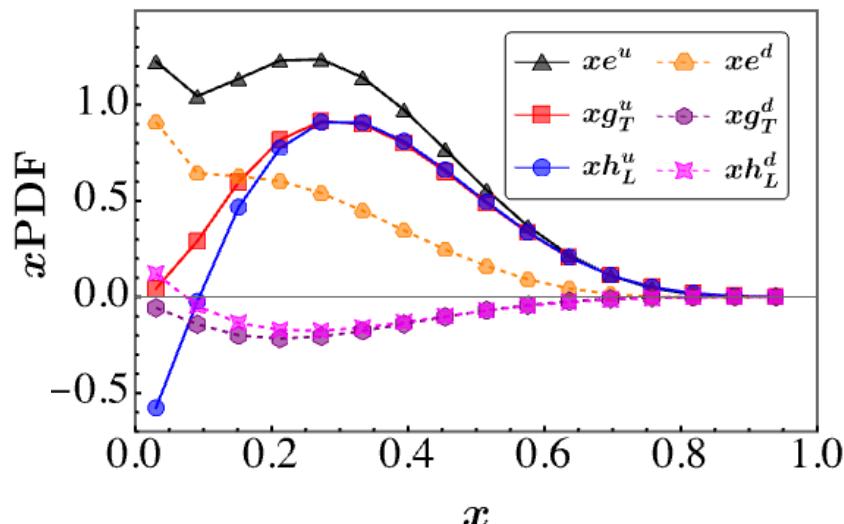
x PDFs of proton

[Zhimin ZHU et al, in preparation]

twist-2 x PDFs

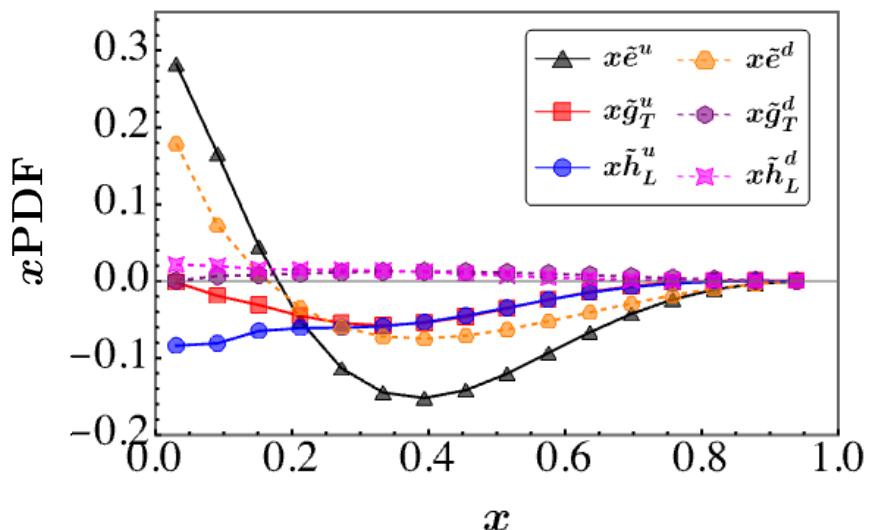


twist-3 x PDFs



genuine twist-3 x PDFs

$$\int \frac{d^2 k_\perp}{(2\pi)^2} f(x, k_\perp) = f(x)$$

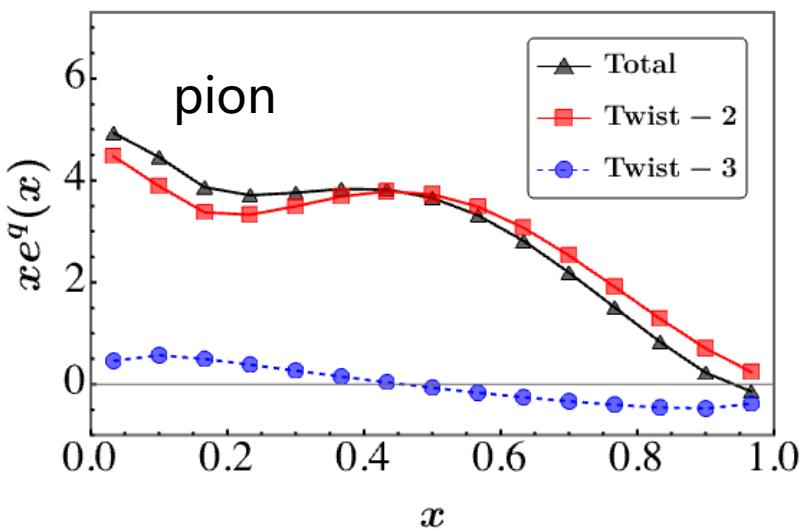
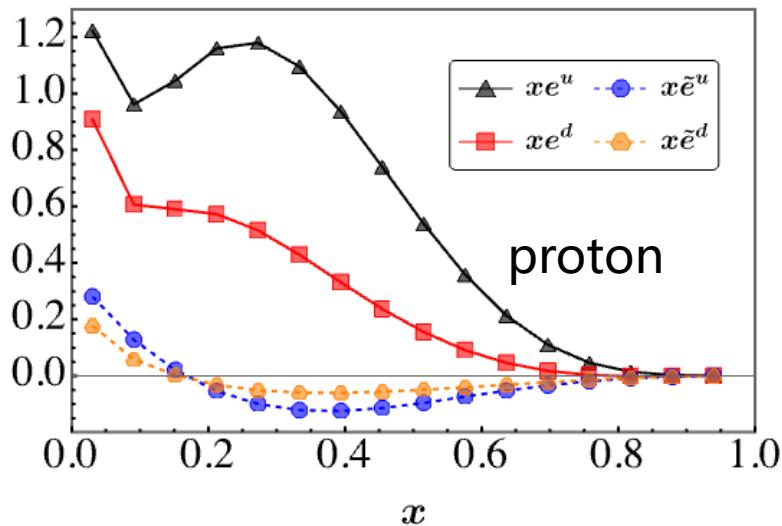


Twist-3 PDFs:

- more concentrating in small x
- larger than twist-2 parts

Twist-3 x PDFs of proton & pion

Sum rules



$$1. \quad \int dx x e^q(x) = \frac{m_q}{M} N_q$$

$$\int dx f_1^q(x) = N_q$$

$$2. \quad e(x, k_\perp) = \frac{m}{M} \frac{f_1(x, k_\perp)}{x} + \tilde{e}(x, k_\perp),$$

proton:

$$\int dx x \tilde{e}(x) = -0.011 \ll \frac{m_q}{M} = \frac{0.31}{0.98}$$

pion:

$$\int dx x \tilde{e}(x) = -0.010 \ll \frac{m_q}{M} = \frac{0.39}{0.139}$$

genuine twist-3 PDFs satisfy
QCD sum rules

Proton mass decomposition

[X.D. Ji, NPB (2020) 115181]

1. $M = M_q^{\text{LF}} + M_g^{\text{LF}} + M_a^{\text{LF}}$

2. $M_q^{\text{LF}} = (a + b)M/2$

$$M_g^{\text{LF}} = (1 - a)M/2$$

$$M_a^{\text{LF}} = (1 - b)M/2$$

$$a = \sum_q \int dx x f_1^q(x) \quad \text{quark kinetic motions}$$

$$b = \sum_q \sigma_q (m_q/M) \quad \text{scalar charge}$$

3. Without gluons $a \rightarrow 1$

With one gluon

Proton

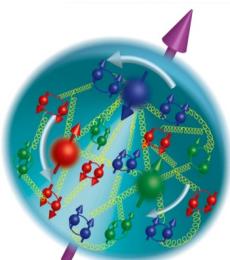
$$a = 0.80$$

$$b = 2.38$$

$$M_q^{\text{LF}} = 1.59 \text{ GeV}$$

$$M_g^{\text{LF}} = 0.10 \text{ GeV}$$

$$M_a^{\text{LF}} = -0.69 \text{ GeV}$$



$$\sigma_q = \int_{-1}^1 dx e_q(x, \mu)$$

Pion

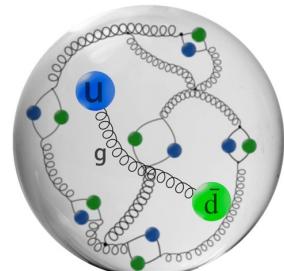
$$a = 0.78$$

$$b = 14.55$$

$$M_q^{\text{LF}} = 1.065 \text{ GeV}$$

$$M_g^{\text{LF}} = 0.0153 \text{ GeV}$$

$$M_a^{\text{LF}} = -0.942 \text{ GeV}$$



Outlook

- TMDs carry a wealth of information about parton structure of hadrons.
- Higher twist TMDs contain more structure information, but are more complex.
- Twist-3 observables need dynamical gluons.
- BLFQ can obtain twist-2 and genuine twist-3 TMDs systematically.
- At initial scale, twist-3 TMDs are larger than twist-2 TMDs.
- Further study: Cross section, T-odd TMDs, twist-3 Qiu-Stermen functions,

Thank You !